

Advanced Topics in Condensed Matter

Lecture 7: Phonons – Anharmonicity Effects (part II)

Dr. Ivan Zaluzhnyy

Prof. Dr. Frank Schreiber



EBERHARD KARLS
UNIVERSITÄT
TÜBINGEN



What to remember (from last week)

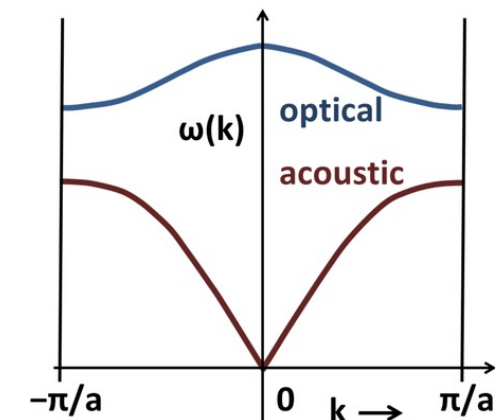
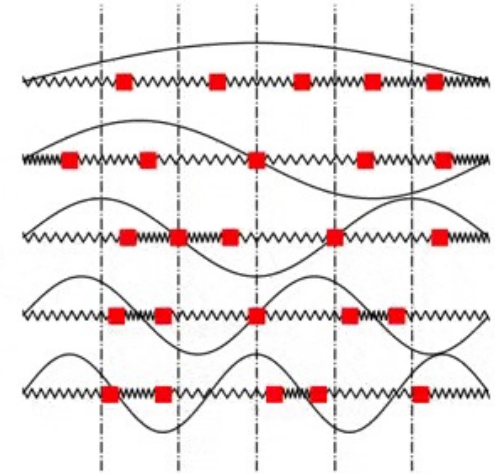
- Elastic scattering is sensitive to the time-averaged structure
- Inelastic scattering is sensitive to the dynamics

$$G(\vec{r}, t) = \frac{1}{N} \left\langle \sum_{n,m=1}^N \delta \left(\vec{r} - ((\vec{r}_n(t) - \vec{r}_m(0))) \right) \right\rangle$$

$$S_{coh}(\vec{q}, \omega) = \frac{1}{2\pi\hbar} \int G(\vec{r}, t) e^{-i\vec{q}\vec{r}} e^{i\omega t} d\vec{r} dt$$

- Inelastic neutron scattering is the best tool to measure phonon dispersion over entire range of ω and k

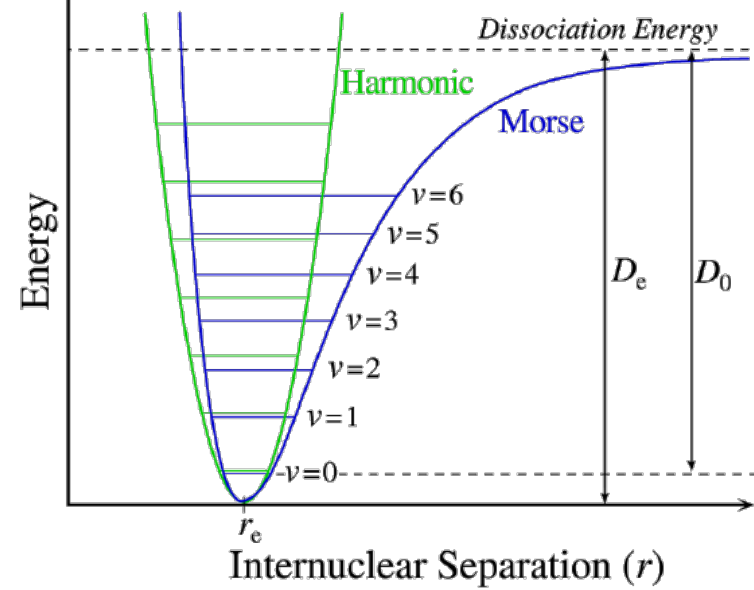
©2017, Bhaskar Kamble



- Raman scattering involves optical phonons (near $k=0$)
- Brillouin scattering involves acoustic phonons (near $k=0$)

Effects related to anharmonicity

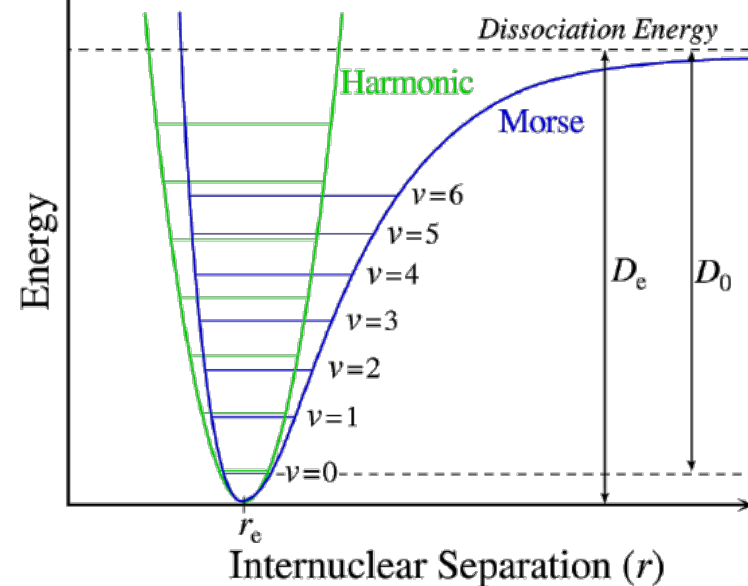
$$V(r) = a(r-r_e)^2 - b(r-r_e)^3 + \dots$$



Effects related to anharmonicity

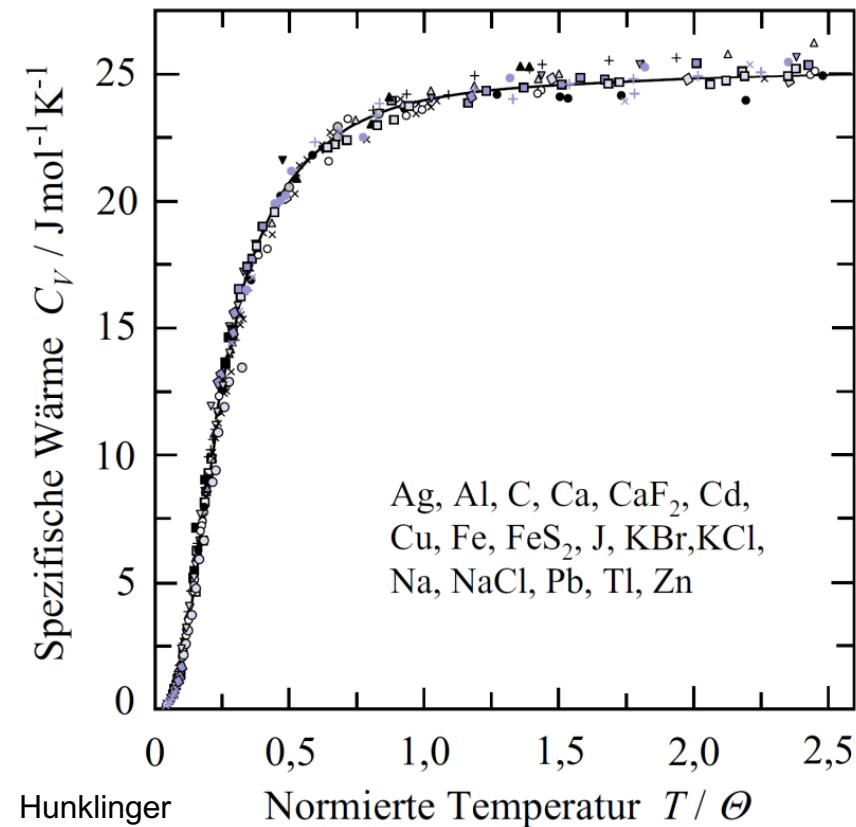
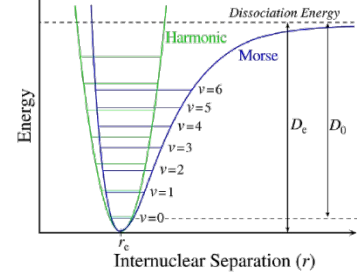
$$V(r) = a(r-r_e)^2 - b(r-r_e)^3 + \dots$$

- the blue color of water
- thermal expansion α of the lattice
- elasticity parameters $C_{(\dots)}$ or B and G are T -dependent
- heat capacity is not strictly constant at high T ;
deviation from Dulong-Petit rule $c = 3Nk_B$
since equipartition theorem demands Hamiltonian $\sim p^2$ and $\sim x^2$
- interaction of phonons with each other (phonon collisions)
- finite lifetime of phonons
wave solutions with linear superposition not stable
with nonlinear terms $\mathcal{O}(u^2)$, i.e. $\partial_t \partial_t u = C \partial_x \partial_x u + \mathcal{O}(u^2)$
- heat conduction by phonons
- Grüneisen parameter connects α with $c(T)$ etc



Fourth:

Thermodynamic implications of anharmonicity

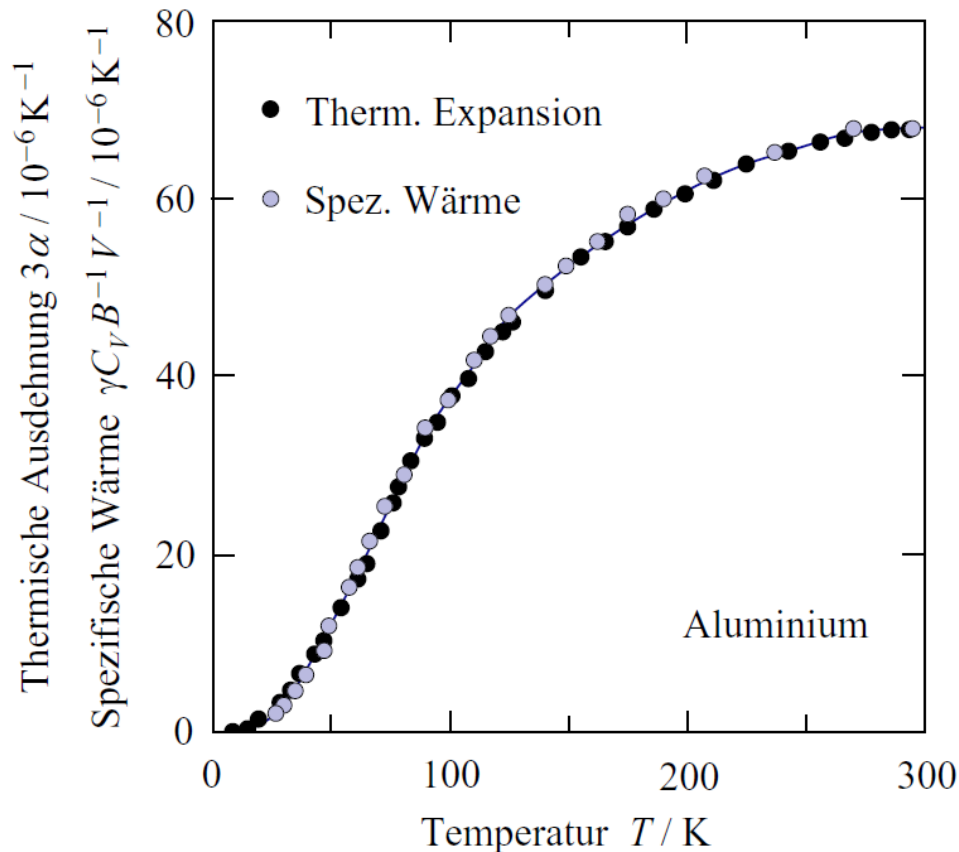
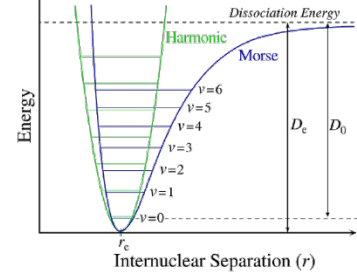


- Violation of the Dulong-Petit law, i.e. heat capacity is *not* constant (and above $3Nk_B$) for high T

Fourth:

Thermodynamic implications

of anharmonicity

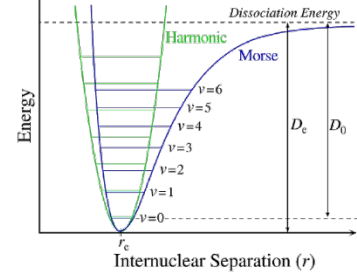


thermal expansion α has very similar T dependence as specific heat $c(T)$; explanation by Grüneisen based on thermodynamics

$$\gamma = (\alpha B) / (c \rho)$$

Fourth:

Thermodynamic implications of anharmonicity



Note the numerous versions
of the Grüneisen parameter
(and the miracles of thermodynamics)

In condensed matter, **Grüneisen parameter** γ is a dimensionless thermodynamic parameter named after German physicist Eduard Grüneisen, whose original definition was formulated in terms of the phonon nonlinearities.^[1]

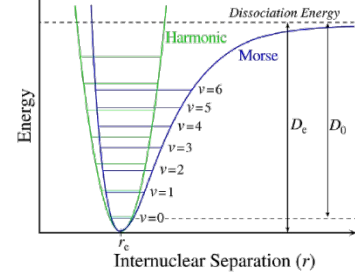
Because of the equivalences of many properties and derivatives within thermodynamics (e.g. see Maxwell relations), there are many formulations of the Grüneisen parameter which are equally valid, leading to numerous interpretations of its meaning. Some formulations for the Grüneisen parameter include:

$$\gamma = V \left(\frac{dP}{dE} \right)_V = \frac{\alpha K_T}{C_V \rho} = \frac{\alpha K_S}{C_P \rho} = \frac{\alpha v_s^2}{C_P} = - \left(\frac{\partial \ln T}{\partial \ln V} \right)_S$$

where V is volume, C_P and C_V are the principal (i.e. per-mass) heat capacities at constant pressure and volume, E is energy, S is entropy, α is the volume thermal expansion coefficient, K_S and K_T are the adiabatic and isothermal bulk moduli, v_s is the speed of sound in the medium, and ρ is density. The Grüneisen parameter is dimensionless.

Fourth:

Thermodynamic implications of anharmonicity



The Grüneisen parameter γ can be related to the anharmonicity.

A proper description of γ is a stringent test for any type of interaction potential.

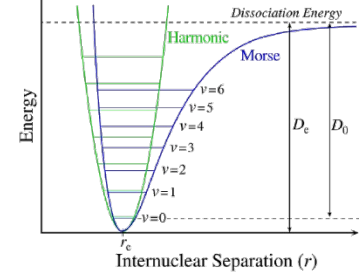
| Lattice | Dimensionality (d) | Lennard-Jones potential | Mie Potential | Morse potential |
|--------------------|------------------------|------------------------------|----------------------------------|--|
| Chain | 1 | $10\frac{1}{2}$ | $\frac{m+n+3}{2}$ | $\frac{3\alpha a}{2}$ |
| Triangular lattice | 2 | 5 | $\frac{m+n+2}{4}$ | $\frac{3\alpha a - 1}{4}$ |
| FCC, BCC | 3 | $\frac{19}{6}$ | $\frac{n+m+1}{6}$ | $\frac{3\alpha a - 2}{6}$ |
| "Hyperlattice" | ∞ | $-\frac{1}{2}$ | $-\frac{1}{2}$ | $-\frac{1}{2}$ |
| General formula | d | $\frac{11}{d} - \frac{1}{2}$ | $\frac{m+n+4}{2d} - \frac{1}{2}$ | $\frac{3\alpha a + 1}{2d} - \frac{1}{2}$ |

Note that analytical calculations are possible for model potentials

Wikipedia

Fourth:

Thermodynamic implications of anharmonicity



The Grüneisen parameter γ can be related to the anharmonicity.

A proper description of γ is a stringent test for any type of interaction potential.

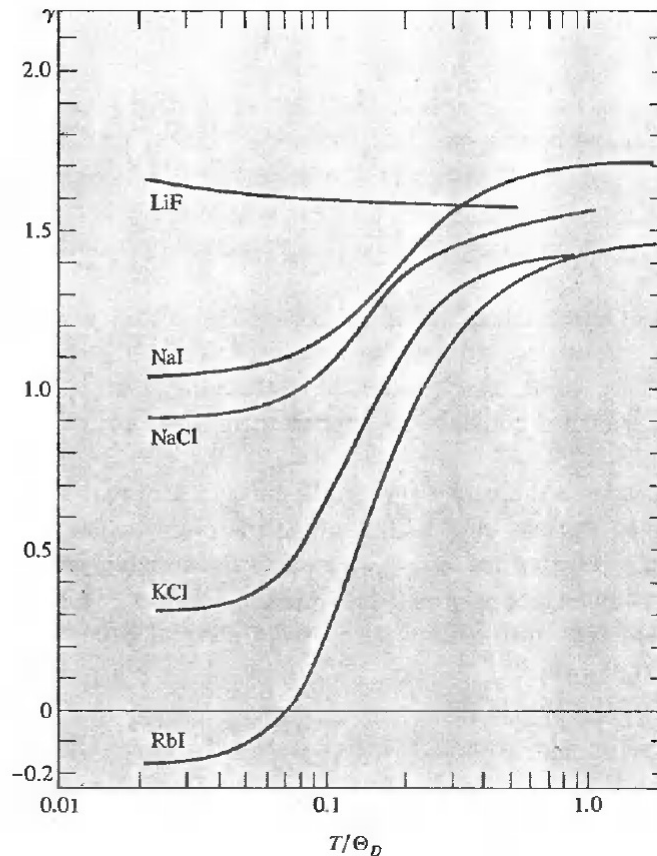


Figure 25.1

Grüneisen parameter vs. T/Θ_D
for some alkali halide crystals.
(From G. K. White, *Proc. Roy Soc. London A286*, 204 (1965).)



What to remember

Harmonic approximation and structural dynamics

- inelastic neutron scattering still best tool to measure $\omega(\mathbf{k})$
- phonon dispersion $\omega(\mathbf{k})$ can be calculated well

Anharmonicity leads to many relevant additional effects

- the blue color of water
- thermal expansion α of the lattice
- elasticity parameters $C_{(\dots)}$ or B and G are $f(T)$
- heat capacity is not strictly constant at high T ;
deviation from Dulong-Petit rule $c = 3Nk_B$
- interaction of phonons with each other (phonon collisions)
- finite lifetime of phonons
wave solutions with linear superposition not stable
with nonlinear terms $\mathcal{O}(u^2)$, i.e. $\partial_t \partial_t u = C \partial_x \partial_x u + \mathcal{O}(u^2)$
- heat conduction by phonons
- Grüneisen parameter connects α with $c(T)$ etc

